

# Stability of a travelling nerve impulse

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## Outline of the talk

1. The model studied and its features. Why Hodgkin-Huxley or FitzHugh-Nagumo?
2. A travelling wave solution: properties and features
3. Stability of a travelling wave: the Evans function and its construction.

## Original references

Hodgkin-Huxley: 1952 **A quantitative description of membrane current** (J. Physiol)

Cole, Antosiewicz, Rabinovitz: **Automatic computation** 1955

FitzHugh, 1961: **Impulses and physiological states in theoretical models of nerve membrane**

Evans and Shank 1970: **Solution to axon equations** (Biophys. J)

Evans 1970-1975: **The stability of the nerve influx** (4 papers Indiana Math. J.)

Hastings 1975: **On travelling wave solutions of the HH eqs.**

# 1 Modeling of the problem

## 1.1 Electrophysiological model

Membrane current density  $I = C_M \partial_t V + I_K + I_{Na} + I_l$ :

Channel currents  $I_K, I_{Na}, I_l$ , Capacitance current  $C_M \partial_t V$  (the membrane being a dielectric with large conductivity), with

$$I_K = \bar{g}_K n^4 (V - V_K), I_{Na} = \bar{g}_{Na} m^3 h (V - V_{Na}), I_l = \bar{g}_l (V - V_l).$$

The quantities  $n, m, h$  represent probabilities of fluxes of ions through the membrane, satisfying:

$$\partial_t n = \lambda_n(V)(n_\infty(V) - n)$$

$$\lambda_n(V) = 0.1 \frac{1 + \frac{V}{10}}{e^{1 + \frac{V}{10}} - 1} + \frac{1}{8} e^{\frac{V}{80}}, (n_\infty(V))^{-1} = 1 + \frac{5}{4} e^{\frac{V}{80}} \frac{e^{1 + \frac{V}{10}} - 1}{1 + \frac{V}{10}}$$

(similar expressions and relations for  $m$  and  $h$ ).

Let  $N_\infty(V) = (n_\infty(V), m_\infty(V), h_\infty(V))^T$

**For**  $(n, m, h) = N_\infty(V)$ ,  $I_K + I_{Na} + I_l = 0$  **iff**  $V = 0$

Numerical values  $\bar{g}_K = 36$ ,  $\bar{g}_{Na} = 120$ ,  $\bar{g}_l = 0.333$ ,

$V_K = 12mV$ ,  $V_{Na} = -115mV$ ,  $V_l = -10.613mV$

$n_\infty(0) = 0.317$ ,  $m_\infty(0) = 0.0529$ ,  $h_\infty(0) = 0.596$

$I_K = -4.36mAcm^{-2}$ ,  $I_{Na} = 1.21mAcm^{-2}$ ,  $I_l = 3.18mAcm^{-2}$

and  $\bar{g}_K(n_\infty(0))^4 = 0.363mmho.cm^{-2}$ ,

$\bar{g}_{Na}(m_\infty(0))^3h_\infty(0) = 0.0105mmho.cm^{-2}$ .

**Gauss law and application on the mean of the potential in the axoplasm**

$$\frac{2}{\sigma r} \frac{\partial^2 V}{\partial x^2} = I.$$

## 1.2 Hodgkin-Huxley equation and PDE system

Introduce  $g(n, m, h) = \bar{g}_K n^4 + \bar{g}_{Na} m^3 h + \bar{g}_l > 0$ ,

$S(n, m, h) = \bar{g}_K n^4 V_K + \bar{g}_{Na} m^3 h V_{Na} + \bar{g}_l V_l$ ,

$$\frac{2}{\sigma r} \frac{\partial^2 V}{\partial x^2} = C_M \partial_t V + g(n, m, h)V - S(n, m, h). \quad (1)$$

Introduce  $N = (n, m, h)$ ,  $\Lambda(V) = \text{Diag}(\lambda_n(V), \lambda_m(V), \lambda_h(V))$ .

System of Hodgkin-Huxley equations ( $X = x\sqrt{\frac{\sigma r}{2}}$ )

$$\begin{cases} \partial_t V - \partial_{X^2}^2 V = -g(n, m, h)V + S(n, m, h) = -g(N)V + S(N) \\ \partial_t N^T = \Lambda(V)(N_\infty(V) - N^T) \end{cases} \quad (2)$$

## 2 Aim of the study of this system

Q1: Stability at rest  $V = 0$ ?

Q2: Existence of progressive solutions  $V(t - \frac{X}{s})$ ?

Q3: Stability of such a progressive solution?

Why study FitzHugh-Nagumo system:

$$\begin{cases} \partial_t V - \partial_{X^2}^2 V = c(y + V - \frac{V^3}{3}) \\ \partial_t y = -c^{-1}(V - a + by) \end{cases} \quad ? \quad (3)$$

- give rise to a Hopf bifurcation when considering the associated system of ODE  $\partial_t V = c(y + V - \frac{V^3}{3}), \partial_t y = -c^{-1}(V - a + by)$  under conditions  $1 - \frac{2b}{3} < a < 1, 0 < b < 1, b < c^2$  for which there is only one equilibrium point which is a stable node.
- be the simplest non trivial nonlinear system appearing when taking into account the behavior of  $\lambda_n(V)$  and  $\lambda_h(V)$  (which are small)  $\Rightarrow$  study of the system

$$\begin{cases} \partial_t V - \partial_{X^2}^2 V = -g(n_\infty(V), m, h_\infty(V))V + S(n_\infty(V), m, h_\infty(V)) \\ \partial_t m = \lambda_m(V)(m_\infty(V) - m) \end{cases} \quad (4)$$

### 3 Study of the rest point $V = 0$

Linearized system at rest:

$$\begin{cases} \partial_t v - \partial_{X^2}^2 v = -g(N_\infty(0))v + \nabla S(N_\infty(0)) \cdot \tilde{n} \\ \partial_n \tilde{n} = \Lambda(0)(N'_\infty(0)v - \tilde{n}) \end{cases}$$

For FitzHugh-Nagumo system one has

$$\begin{cases} \partial_t v - \partial_{X^2}^2 v = c((1 - x_0^2)v + y) \\ \partial_t y = -c^{-1}(bv + y) \end{cases}$$

Exponential stability at rest and construction of the Evans function are consequences of the properties of the matrix

$$A(\tau) = \begin{pmatrix} -g(N_\infty(0)) - \tau & \nabla S(N_\infty(0)) \\ \Lambda(0)N'_\infty(0) & -\Lambda(0) \end{pmatrix}.$$

$$\text{(respectively)} A_f(\tau) = \begin{pmatrix} c(1 - x_0^2) - \tau & c \\ -\frac{b}{c} & -\frac{1}{c} \end{pmatrix}.$$

Properties  $A = A(0) = \begin{pmatrix} a & r \\ c & -\Lambda(0) \end{pmatrix}$  has four eigenvalues of

negative real part. Moreover  $A(\tau)$  has no imaginary eigenvalues for  $\tau \geq 0$ .

The eigenvalues of  $A(\tau)$  are solution of

$$R(\lambda) - \tau = -g(N_\infty(0)) - \tau - \lambda + \nabla S(N_\infty(0))(\Lambda(0) + \lambda I)^{-1} \Lambda(0)N'_\infty(0) = 0.$$

Note that for  $A_f(\tau)$  the result is similar.

**Theorem 1** *The system (2) is exponentially stable at rest iff all eigenvalues of  $A(\xi^2)$  stay in the half plane  $\Re\lambda \leq -\kappa < 0$ . It is equivalent to the exponential stability of the linearized system.*

This is the consequence of the properties of the eigenvalues and of the representation of the solution of the system as

$$W_\xi(t) = F \star D + \begin{pmatrix} 0 & 0 \\ 0 & e^{-t\Lambda} \end{pmatrix} D, D = W_\xi(0),$$

where  $F(\xi, t) = e^{tA(\xi^2)} - \begin{pmatrix} 0 & 0 \\ 0 & e^{-t\Lambda} \end{pmatrix}$  is exponentially controlled,

hence

$$\|W_\xi(t)\| \leq C \|W_\xi(0)\| e^{-\kappa t}.$$

## 4 Travelling wave construction

Q2: Travelling wave  $(V_0(t - \frac{X}{s}), n_0(t - \frac{X}{s}), m_0(t - \frac{X}{s}), h_0(t - \frac{X}{s}))$ .

Dynamical system

$$\begin{cases} V_0' - \frac{1}{s^2}V_0'' = -g(N_0)V_0 + S(N_0) \\ N_0' = \Lambda(V_0)(N_\infty(V_0) - N_0) \end{cases} \quad (5)$$

It is equivalent to

$$\frac{d}{d\tau} \begin{pmatrix} V_0 \\ W_0 \\ N_0 \end{pmatrix} = \begin{pmatrix} s^2W_0 \\ -g(N_0)V_0 + S(N_0) + s^2W_0 \\ \Lambda_0(V_0)(N_\infty(V_0) - N_0) \end{pmatrix} \quad (6)$$

Linearized system around  $(0, 0, n_\infty(0), m_\infty(0), h_\infty(0))^t$ ,

$U = (\tilde{V}, s^{-2}\tilde{V}', \tilde{n}, \tilde{m}, \tilde{h})^t$ :

$$U' = \begin{pmatrix} 0 & s^2 & 0 & 0 & 0 \\ -a & s^2 & -r_1 & -r_2 & -r_3 \\ c_1 & 0 & -\lambda_1 & 0 & 0 \\ c_2 & 0 & 0 & -\lambda_2 & 0 \\ c_3 & 0 & 0 & 0 & -\lambda_3 \end{pmatrix} U = M(s^2)U.$$

The eigenvalues of this matrix are solution of  $R(\lambda) = \frac{\lambda^2}{s^2}$ , that is  $\lambda \in \text{Sp}(A(-\frac{\lambda^2}{s^2}))$ . No imaginary eigenvalue. Only one eigenvalue  $\lambda_0(s^2) > s^2$  in  $\Re > 0 \Rightarrow$  unstable manifold of dimension 1, stable manifold of dimension 4.

## Construction of the unstable wave

**Theorem 2** *For all  $s \neq 0$  there exist two solutions  $\phi_+(\tau)$  et  $\phi_-(\tau)$  such that all solutions of (5) which go to 0 at  $-\infty$  are  $U = \phi_+(\tau - \tau_0)$  or  $U = \phi_-(\tau - \tau_0)$ .*

Constructive proof: The eigenvector  $E_+$  associated with  $\lambda_0(s^2)$  is

$$E_+ = \left(1, \frac{s^2}{\lambda_0(s^2)}, \frac{c_1}{\lambda_1 + \lambda_0(s^2)}, \frac{c_2}{\lambda_2 + \lambda_0(s^2)}, \frac{c_3}{\lambda_3 + \lambda_0(s^2)}\right)^t$$

As (2) writes  $U' = F'(U_\infty)(U - U_\infty) + R(U - U_\infty)$  construct a

Picard sequence: 
$$\begin{cases} Y_{n+1}^K(\tau) = U_\infty + \int_{-\infty}^{\tau} R(Y_n(\tau') - U_\infty) d\tau' \\ Y_0^K(\tau) = U_\infty + K E_+ e^{\lambda_0(s^2)\tau} \end{cases} . \text{ As}$$

$Y^K(\tau) = Y^1\left(\tau - \frac{1}{\lambda_+(s^2)} \ln |K|\right)$ , we define the two solutions associated with  $K > 0$  and  $K < 0$ .

## Finding a travel speed

**A travelling wave solution is in the intersection of the unstable manifold and the stable manifold.**

- Stable manifold: dimension 4: implicit equation

$$\Phi_{s^2}(V, W, n, m, h) = 0.$$

Condition for the travel speed:  $\Phi_{s^2}(\phi_+(\tau)) \cdot \Phi_{s^2}(\phi_-(\tau)) = 0$ .

- Numerics: impossible as an IVP: stiff (large eigenvalues) problem.

Possible as a BVP: non zero solution satisfying the four conditions at  $\tau = -T$  and one condition at  $\tau = T$  (not yet done)

Remark: it is the solution of a minimization problem coupled with the ODEs (work in progress):

$$s^{-2}V'' = V' + g(N)V - S(N) \Leftrightarrow -s^{-2}u'' + (g(N) + \frac{s^2}{4})u = S(N)e^{-\frac{s^2}{2}t}$$

on  $H^1([T_0, +\infty[)$ . Coercivity constant  $\alpha = \min(s^{-2}, g_l + \frac{s^2}{4})$ .

### Construction of the stable manifold:

Let  $P_{s^2}$  be the matrix of eigenvectors of  $M(s^2)$  (in the order positive e.v. and e.v. of negative real part). Let us write  $F = P_{s^2}^{-1}U$ . The dynamical system writes  $F' = D(s^2)F + g(F, s^2)$ ,  $F$  solution of

$$F(t) = U_1(t)(a_1, a_2, a_3, a_4, 0)^T + \int_0^t U_1(t-t')g(F(t'), s^2)dt' - \int_t^{+\infty} (0, 0, 0, 0, g_5(F(t'), s^2))^T e^{\lambda_0(s^2)(t-t')} dt'.$$

This functional equation has a unique solution when initial value  $(a_1, a_2, a_3, a_4, 0)^T$  is given. One has  $F_5(0, a) = - \int_0^\infty g_5(F(t', a), s^2)e^{\lambda_0(s^2)(t-t')} dt'$  (equation of the stable manifold).

**Q3: Stability of solutions of**

$$\begin{cases} v'_0 + s^{-2}v''_0 + (-g(n_0)v_0 + S(n_0)) = 0 \\ n'_0 + \Lambda_\infty(v_0)(n_\infty(v_0) - n_0) = 0 \end{cases} \quad ?$$

System in the adapted variables:  $\xi = \frac{x}{s} - t, t$

$$\begin{cases} \partial_\xi v + s^{-2}\partial_{x^2}^2 v + [-g(n)v + S(n)] = \partial_t v \\ \partial_\xi n + \Lambda_\infty(v)(n_\infty(v) - n) = \partial_t n \end{cases} \quad (7)$$

If one derivates the system on  $v_0, n_0$  one gets

$$\begin{cases} v''_0 + s^{-2}v'''_0 - g(n_0)v'_0 + (-g'(n_0)v_0 + S'(n_0))n'_0 = 0 \\ n''_0 + [\Lambda'_\infty(v_0)(n_\infty(v_0) - n_0) + \Lambda_\infty(v_0)n'_\infty(v_0)]v'_0 - \Lambda_\infty(v_0)n'_0 = 0. \end{cases}$$

This system writes

$$\begin{pmatrix} v'_0 \\ n'_0 \end{pmatrix}' + s^{-2} \begin{pmatrix} v'_0 \\ 0 \end{pmatrix}'' + B(\xi) \begin{pmatrix} v'_0 \\ n'_0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \text{ where}$$

$$B(\xi) = A(0) + P(\xi), P(\xi) \rightarrow 0 \text{ for } |\xi| \rightarrow +\infty.$$

The linearization of (7) is

$$T_0 \begin{pmatrix} \tilde{v} \\ \tilde{n} \end{pmatrix} = \begin{pmatrix} \tilde{v} \\ \tilde{n} \end{pmatrix}' + s^{-2} \begin{pmatrix} \tilde{v} \\ 0 \end{pmatrix}'' + B(y) \begin{pmatrix} \tilde{v} \\ \tilde{n} \end{pmatrix} = \partial_t \begin{pmatrix} \tilde{v} \\ \tilde{n} \end{pmatrix}.$$

Spectrum of operator  $T_0$  gives the stability of the travelling wave.

Introduce  $T_0^a$  the adjoint operator of  $T_0$ . The kernel of  $T_0^a$  is of dimension 1 (because  $(v'_0, n'_0)$  is in the kernel of  $T_0$  and the dimension is 1). Denote by  $\gamma(\xi)$  a function in  $Ker T_0^a$ .

Evans showed that  $\int \gamma(\xi) \cdot (v'_0(\xi), n'_0(\xi)) d\xi = 0$  is equivalent to the existence of a  $\psi$  solution of  $T_0\psi = (v'_0(\xi), n'_0(\xi))^T$ .

Main result:

The system is exponentially stable at  $\begin{pmatrix} v'_0(\xi) \\ n'_0(\xi) \end{pmatrix}$  iff

$\int \gamma(\xi) \cdot \begin{pmatrix} v'_0(\xi) \\ n'_0(\xi) \end{pmatrix} d\xi \neq 0$  and the maximum of the real parts of the

$\lambda \neq 0$  such that  $T_0\psi = \lambda\psi$  has a bounded solution is negative.

## 5 Evans function construction and formulation

### 5.1 Construction of the bounded solutions at $\pm\infty$ of the linearized system

$$\frac{d}{d\xi} \begin{bmatrix} \tilde{V} \\ \tilde{W} \\ \tilde{N} \end{bmatrix} + \begin{pmatrix} 0 & s^2 & 0 \\ a(\xi) - \lambda & -s^2 & r(\xi) \\ c(\xi) & 0 & -\Lambda(\xi) - \lambda I \end{pmatrix} \begin{bmatrix} \tilde{V} \\ \tilde{W} \\ \tilde{N} \end{bmatrix} = 0 \quad (8)$$

$$-\frac{d}{d\xi} \begin{bmatrix} \tilde{V}^* \\ \tilde{S}^* \\ \tilde{N}^* \end{bmatrix} + \begin{pmatrix} 0 & a(\xi) - \lambda & c^T(\xi) \\ s^2 & -s^2 & 0 \\ 0 & r^T(\xi) & -\Lambda(\xi) - \lambda I \end{pmatrix} \begin{bmatrix} \tilde{V}^* \\ \tilde{S}^* \\ \tilde{N}^* \end{bmatrix} = 0 \quad (9)$$

Uniqueness of the eigenvalue  $\nu(\lambda)$  of positive real part of

$$\begin{pmatrix} 0 & s^2 & 0 \\ a - \lambda & -s^2 & r \\ c & 0 & -\Lambda - \lambda I \end{pmatrix}, \Re \lambda > \delta_0 > 0, \text{ eigenvector}$$

$$B_+ = \left(1, \frac{\nu(\lambda)}{s^2}, (\Lambda + (\lambda + \nu(\lambda))I)^{-1}c\right)^T$$

and the eigenvector for the adjoint matrix

$$B_- = \left(1 + \frac{\nu(\lambda)}{s^2}, 1, (\Lambda + (\lambda + \nu(\lambda))I)^{-1}r^T\right)^T$$

Construction of the associated solution at  $-\infty$  for the direct system (8) and the adjoint system (9) with a prescribed behavior at infinity:

$$\|\beta(\lambda, \xi) - e^{\nu(\lambda)\xi} B_+\| + \|\beta^*(\lambda, -\xi) - e^{\nu(\lambda)\xi} B_-\| \leq C e^{(\nu(\lambda) + \epsilon_0)\xi}$$

## 5.2 Definition of the Evans function

$$D(\lambda) = \beta(\lambda, \xi) \cdot \beta^*(\lambda, \xi), \forall \xi, \text{ independant of } \xi. \quad (10)$$

Properties:

**Theorem 3** a) *The function  $D(\lambda)$  is an analytic function, such that  $D(0) = 0$ .*

b) *The system (2) is exponentially stable for the profile  $V_0$  with speed  $s_0$  iff  $D(\lambda) \neq 0$  for  $\lambda, \Re \lambda \geq 0, \lambda \neq 0$  and  $\frac{\partial D}{\partial \lambda}(0) \neq 0$ .*

c) *If  $\frac{\partial D}{\partial \lambda}(0) > 0$ , no stability. If  $\frac{\partial D}{\partial \lambda}(0) < 0$ , one cannot conclude.*

d) *The sign of  $\frac{\partial D}{\partial \lambda}(0) > 0$  is the sign of*

$$\lim_{\xi \rightarrow -\infty} e^{\lambda_+(s_0^2)\xi} \begin{pmatrix} 1 + \frac{\lambda_+(s_0^2)}{s_0^2} \\ 1 \\ (\Lambda + (\lambda_+(s_0^2))I)^{-1} r^T \end{pmatrix} \partial_s U(s_0, \xi).$$

Note:  $\lambda$  root of order  $m$  of  $D$  iff  $\lambda$  eigenvalue of multiplicity  $m$  of  $T_0$ .

## 6 Conclusion and next steps

- Some physiological conclusions easy from a mathematical point of view (reversibility, invariance by translation, generation of this spontaneous influx with a small adapted disturbance)
- Numeric calculations of the 50' not convincing (cannot be reproduced easily with the methods of HH)
- Work in progress from a theoretical point of view and a numerical point of view
- Powerfulness of the Evans function and of the analysis of the stable manifold
- One feature of the Hodgkin-Huxley system has been forgotten: the equation of diffusion is linear in  $v$ .
- What about Evans formulation for FitzHugh-Nagumo?